Physics Letters B ••• (••••) •••–•••

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Physics Letters B

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Corrigendum

Corrigendum to “Determination of the superstring moduli from the higher-derivative terms” [Phys. Lett. B 495 (3–4) (2000) 401–406]

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A R T I C L E I N F O

Article history:
Received 29 August 2015
Accepted 2 September 2015
Available online xxxx
Editor: M. Cvetič

Page	Line	Eq.	Printed as	Should read
401	11		$B_{\mathrm{r}} \equiv \mathrm{e}^{\sqrt{2}\kappa_0\sigma_A}$	$B_{\mathrm{r}} \equiv \mathrm{e}^{\sqrt{2/3}\kappa_0\sigma_B}$
401	41		$\Lambda'^3/M_{\mathrm{P}}^3$	$\Lambda'^3/M_{\mathrm{P}}^2$
402	12	(2)	$+\, B_{\mathrm{r}}^{-1} ds^{-2}$	$+\, B_{\mathrm{r}} d\bar{s}^2$
402	15	(3)	$\left[1 + \frac{15\zeta(3)\bar{\chi}}{16\lambda B_{\mathrm{r}}^3}\right]$	$\left[1 - \frac{15\zeta(3)\bar{\chi}}{16\lambda B_{\mathrm{r}}^3}\right]$
402	16	(3)	$+\, \frac{1}{2} B_{\mathrm{r}}^{-2} (\nabla\theta_A)^2$	$+\, \frac{1}{2} B_{\mathrm{r}}^{-2} (\nabla\theta_B)^2$
402	17	(3)	$-\, \mathrm{Tr}(F\widetilde{F})/30$	$-\, \mathrm{Tr}(FF)/30$
402	18	(3)	$-\, A_{\mathrm{r}}^3 \mathcal{R}^4$	$+\, A_{\mathrm{r}}^3 \mathcal{R}^4$
402	39		$\mathcal{R}$	$\mathcal{R}^2$
402	46	(8)	$\left[1 + \frac{15\zeta(3)\bar{\chi}}{16\lambda B_{\mathrm{r}}^3}\right]^{-1}$	$\left[1 - \frac{15\zeta(3)\bar{\chi}}{16\lambda B_{\mathrm{r}}^3}\right]^{-1}$
403	5		$\bar{\chi} = -6$	$\bar{\chi} = \pm 6$, taking into account mirror symmetry – see P.S. Aspinwall and C.A. Lütkin, Nucl. Phys. B 353 (1991) 427 –
403	7	(13)	$\left(1 - 6.676/\lambda B_{\mathrm{r}}^3\right)^{-1}$	$\left(1 \mp 6.761/\lambda B_{\mathrm{r}}^3\right)^{-1}$
403	13		remains positive	remains positive if $\bar{\chi} > 0$
403	34		$G_{\mathrm{N}}\rho \ll 1$	$G_{\mathrm{N}}^2\rho \ll 1$
403	80	(16)	$V(\sigma_B) = \frac{15}{32} \ldots$	$V(\sigma_B) = -\frac{15}{32} \ldots$
403	81	(16)	$e^{\sqrt{8/3}\kappa_0\sigma_B}$	$e^{-\sqrt{8/3}\kappa_0\sigma_B}$
404	11	(20)	$-\, A_1(t) e^{\sqrt{6}\kappa_0\sigma_B}$	$-\, A_1(t)\,(\bar{\chi}/\lvert\bar{\chi}\rvert)\,\mathrm{e}^{-\sqrt{6}\kappa_0\sigma_B}$
404	11	(20)	$+\, A_2(t) e^{\sqrt{8/3}\kappa_0\sigma_B}$	$+\, A_2(t) e^{-\sqrt{8/3}\kappa_0\sigma_B}$
404	14	(21)	$A_2 = 32 A_{\mathrm{r}} \widetilde{B}(2-\gamma)(\gamma-1)/\gamma^4 t^4$	$A_2 = 32 A_{\mathrm{r}} \widetilde{B}(2-\gamma)(\gamma-1)/3\gamma^4 t^4$
404	15		$\bar{\chi} = -6$	$\bar{\chi} = \pm 6$
404	26–28		If this is the case, then $V(\sigma_B)$ is unbounded from below as $\sigma_B \to \infty$ ($B_{\mathrm{r}} \to 0$) and tends to zero as $\sigma_B \to +\infty$ ($B_{\mathrm{r}} \to 0$), with a maximum value	If this is the case, and if $\bar{\chi} = +6$, then $V(\sigma_B)$ is unbounded from below as $\sigma_B \to -\infty$ ($B_{\mathrm{r}} \to 0$) and tends to zero as $\sigma_B \to +\infty$ ($B_{\mathrm{r}} \to \infty$), with a maximum value
404	29	(22)	$V_{\max} = 4A_2^3/27A_1^2 \approx 10^6 \ldots$	$V_{\max} = 4A_2^3/27A_1^2 \approx 5 \times 10^5 \ldots$
404	31	(23)	$B_{\mathrm{r}} = 3A_1/2A_2 \approx 2 \times 10^{-3} \ldots$	$B_{\mathrm{r}} = 3A_1/2A_2 \approx 6 \times 10^{-3} \ldots$

(*continued on next page*)

DOI of original article: http://dx.doi.org/10.1016/S0370-2693(00)01251-X.
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http://dx.doi.org/10.1016/j.physletb.2015.09.007

(*continued*)

Page	Line	Eq.	Printed as	Should read
404	38		… (see Table 2).	… (see Table 2). If, on the other hand, $\bar{\chi} = -6$, then $V(\sigma_B) \to +\infty$ as $\sigma_B \to -\infty$ ($B_r \to 0$), decreasing monotonically to zero as $\sigma_B(B_r)$ increases to $+\infty$.
404	42		$\Gamma_1(\gamma)$	$\Gamma_2(\gamma)$
404	45		$\frac{4}{57}$ $\frac{4}{3}$ $\frac{52}{21}$ $\frac{4}{3}$ $\frac{188}{87}$ $\frac{52}{51}$	$\frac{10}{21}$ $\frac{4}{3}$ $\frac{10}{3}$ $\frac{4}{3}$ $\frac{38}{15}$ $\frac{10}{3}$
404	49		ξ_ρ	$\xi\rho$
405	22	(25)	$\frac{1}{4}\mathrm{Tr}\left(F\widetilde{F}\right)/30$	$\frac{1}{4}\mathrm{Tr}\,(FF)/30$
405	73		From the formulas (21) and (22),	From the formulae (21) and (22), when $\bar{\chi} > 0$
406	6	(34)	$1.89\lambda^{1/3}$	$1.89/\lambda^{1/3}$
406			Ref. 5 should read [5] M.D. Pollock, Phys. Lett. B 411 (1997) 68; Phys. Lett. B 743 (2015) 542 (C).	

Please cite this article in press as: M.D. Pollock, Corrigendum to "Determination of the superstring moduli from the higher-derivative terms" [Phys. Lett. B 495 (3–4) (2000) 401–406], Phys. Lett. B (2015), http://dx.doi.org/10.1016/j.physletb.2015.09.007